The Electric Field of the Sun and Solar Wind C. Fred Driscoll, UCSD Physics

DPP2020 / TO16.7

Electric fields are endemic to plasmas, and are often concentrated in Sheaths, resulting in particle acceleration.

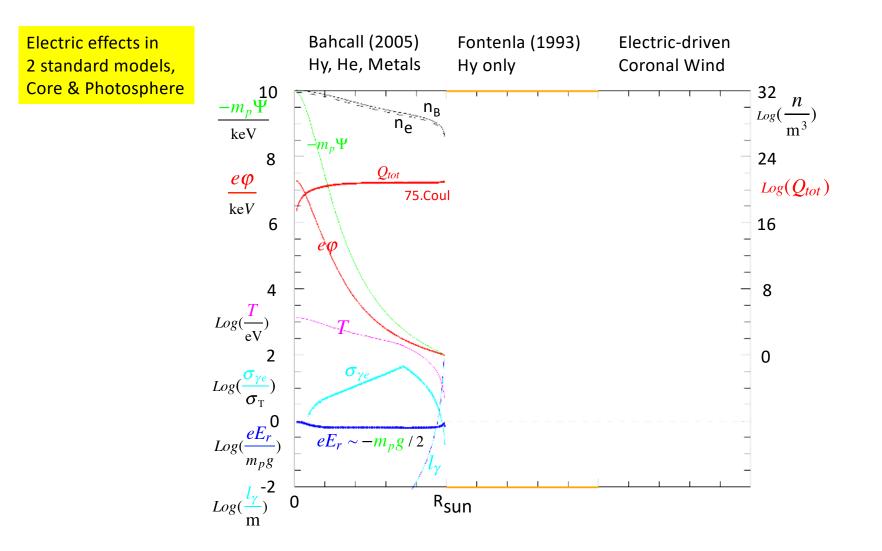
Here, electric effects are calculated for the Solar interior, photosphere, and corona, based on standard 1-D radial models by Bahcall (2005) and Fontenla (1993). The major uncertainty is in the photon-electron scattering cross-section  $\sigma_{\gamma e}$  for re-combining (i.e. correlated) plasmas.

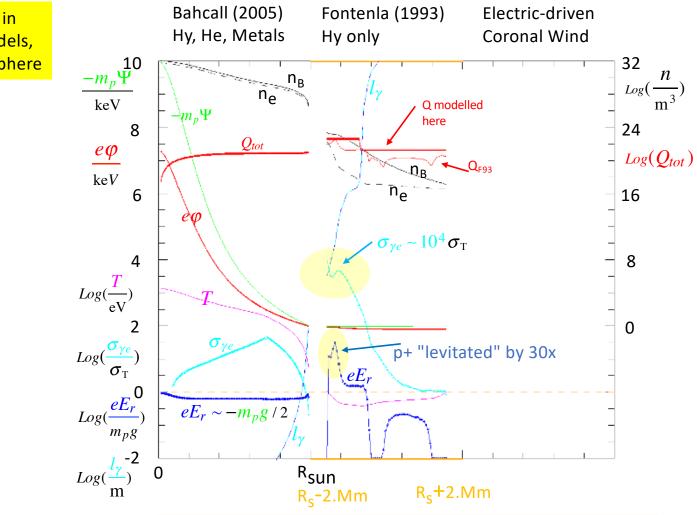
A reasonable estimate for  $\sigma_{\gamma e}$  shows that electric fields in the photospheric sheath and corona can accelerate protons out of the 2keV gravity well and up to the 1.3 keV energies observed in the Solar Wind. The light scattering from the accelerating proton/electron wind approximates that attributed to a hot, hydrostatic K-Corona. Energetically, this requires about 10<sup>-6</sup> of the thermal photon flux of the sun.

The 1-D radial model identifies the energetics as thermo-electric and photo-electric; but a more realistic model would include surface granulation on the 0.5Mm scale, and would probably show filamentary beams with diameters down to the 10.km scale. The inevitable small charge imbalances would generate strong magnetic fields, and filament clumping would increase the visibility and impact of the small beams.

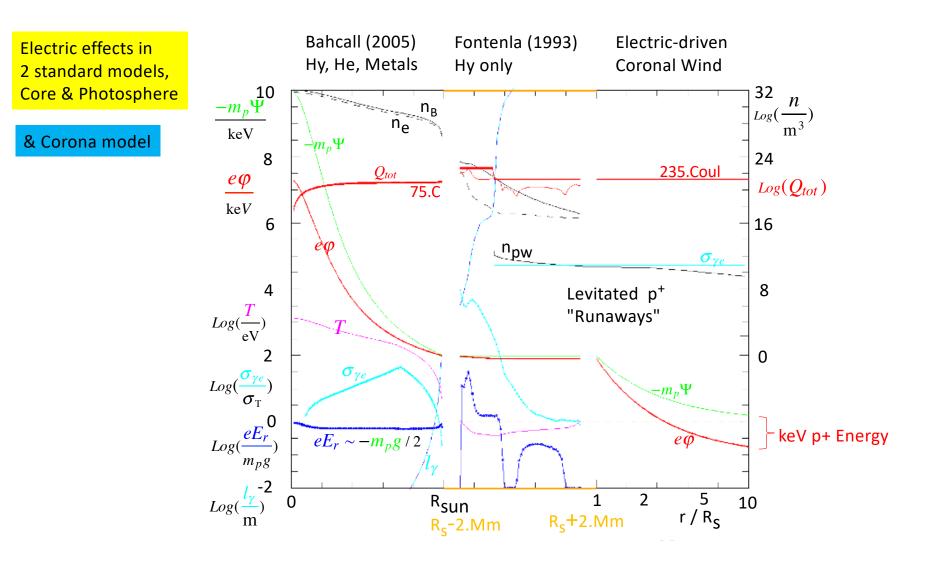
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	Equlibrium Stellar Fluid Eqns:	mass m <sub>p</sub> m <sub>e</sub>	charge e⁻p⁺	photons γ
	$\nabla^2 \Psi(r) = G m_p n_p(r)$	Gravity		
2	$\nabla \cdot \Gamma_{\varepsilon}(r) = \frac{d}{dt} \varepsilon(r) - (4aT^3) T'(r) l_{\gamma} = \frac{4}{c} \Gamma_{\varepsilon}$	Fusion I	Energy Flu	IX
3	$-(4aT^3)T'(r) l_{\gamma} = \frac{4}{c}\Gamma_{\varepsilon}$	Therma	l Energy D	oiffusion
4 <i>a</i>	$[n_pT]' + n_p m_p \Psi' + (+e)n_p \Phi' = 0$		Fluid Mo	
4 <i>b</i>	Thermo- Photo- Electric $[n_eT]' - \frac{\Gamma_{\epsilon\gamma}}{c \ l_{\gamma e}} + n_e m_e \Psi' + (-e)n_e \Phi' = 0$	Electro	on Fluid M	lomentum
4 <i>a</i> -	$\frac{1}{4b} \left[ (2n)T \right]' - \frac{\Gamma_{\varepsilon\gamma}}{c  l_{\gamma e}} + n  m_p \Psi'$	= 0 Tot	al Fluid N	lomentum
4 <i>a</i> –	$4b \qquad \qquad \frac{\Gamma_{\epsilon\gamma}}{c  l_{\gamma e} n_e} +  m_p \Psi' +  (2e)$	$\Phi' = 0$ Ele	ctric Field	ł
	$-m_p g(r)/2$	$\approx e E_{\mathrm{Th}}(r)$	Theri in hig collis	mo-Electric gh-density ional regime
	$\sigma_{\gamma e} \equiv \frac{1}{l_{\gamma e} n_{e}} \qquad \qquad \frac{\Gamma_{\varepsilon \gamma}}{2 c} \sigma_{\gamma e} =$	$eE_{\gamma}$	Photo	o-Electric : Ph T

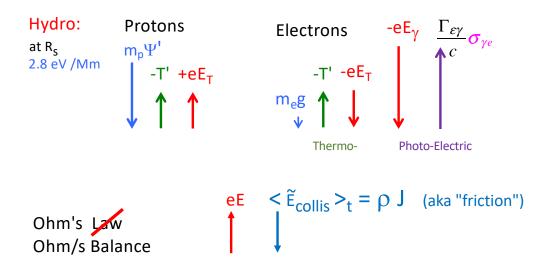




Electric effects in 2 standard models, Core & Photosphere



## Force Balance with Electric Fields



Photon-electron scattering cross-section  $\sigma_{\gamma e}$  increases with plasma density & correlation.

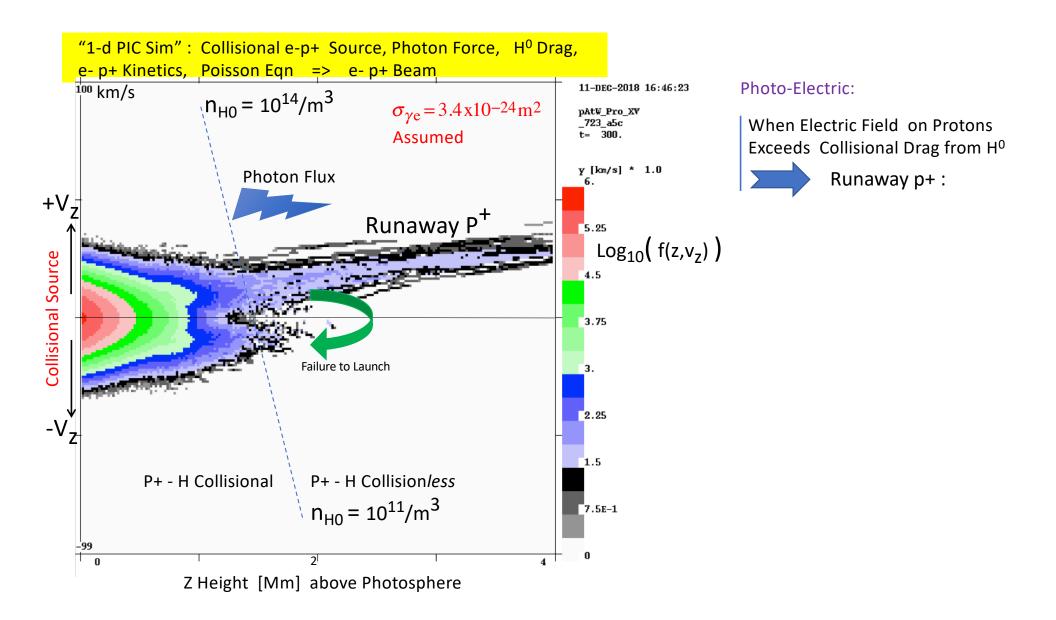
e- p+ Strongly Corelated  $\sigma(\gamma \rightarrow H^*) \sim \pi a_0^2 = 0.6 \text{ x} 10^{-20} \text{ m}^2$   $\sigma(\gamma \rightarrow H^- \quad bf) \sim 0.5 \text{ x} 10^{-20} \text{ m}^2$   $\sigma(\gamma \rightarrow H + e^- ff) \sim 0.5 \text{ x} 10^{-20} \text{ m}^2$ 

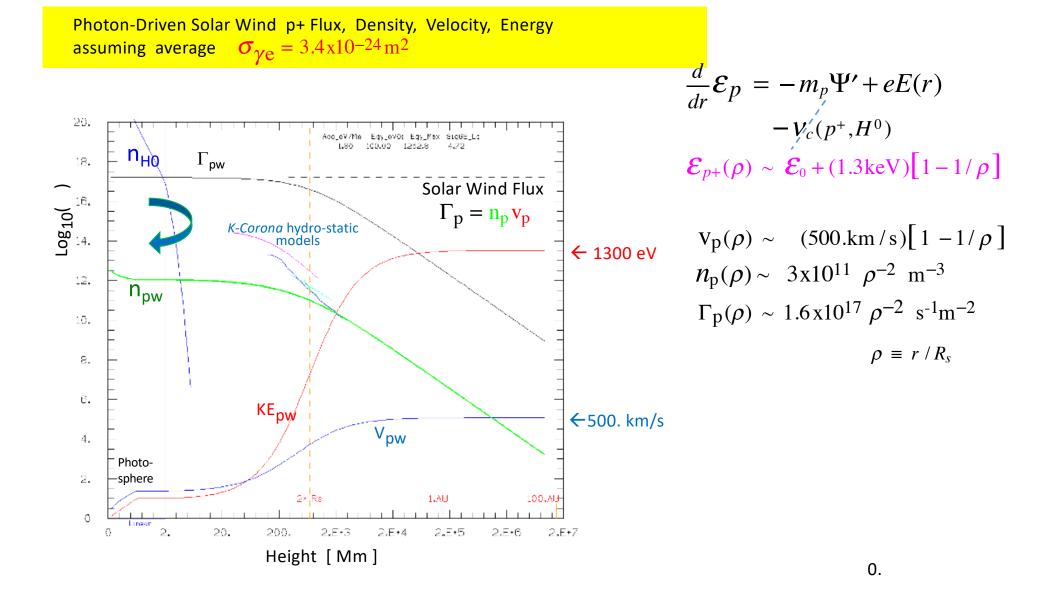
"inverse Bremstrahling" @ n=10<sup>22</sup>, T=1eV

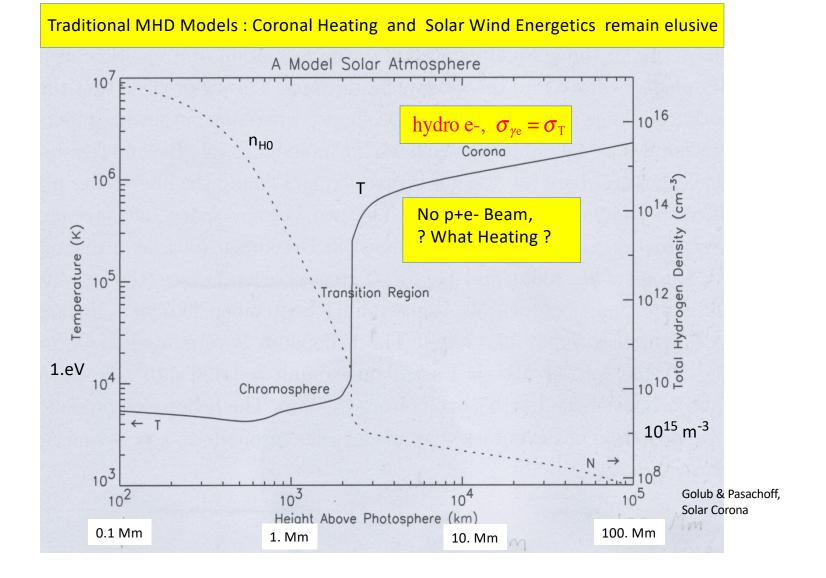
 $\sigma_{\gamma e} \sim 3.4 \, \mathrm{x10^{-24}} = Model$ 

isolated electron

$$\sigma_{\rm T} = 0.7 \, {\rm x} 10^{-28} \, {\rm m}^2$$







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